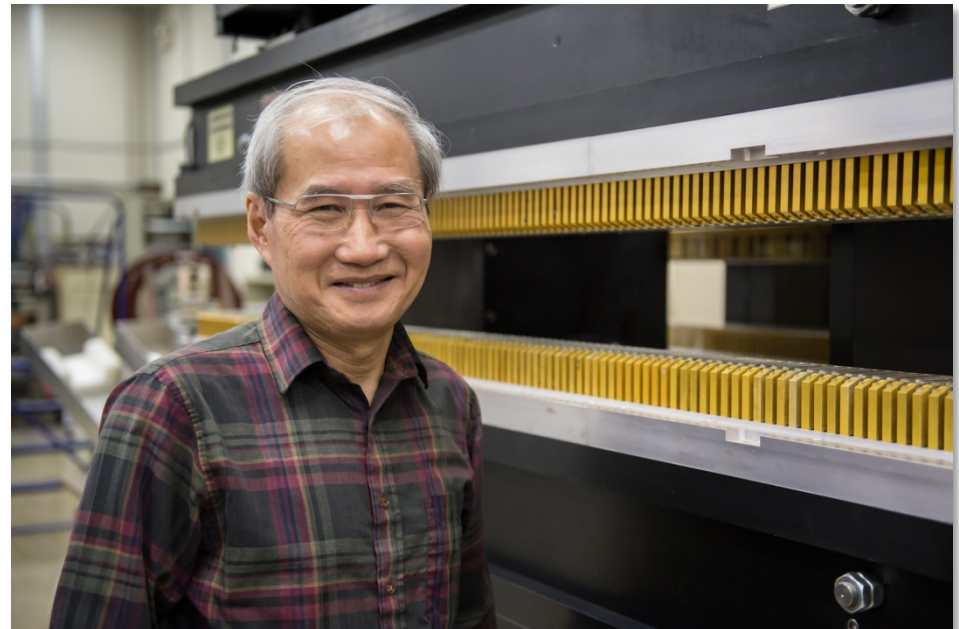


FEL theory: From LEUTL to LCLS

Zhirong Huang (SLAC)

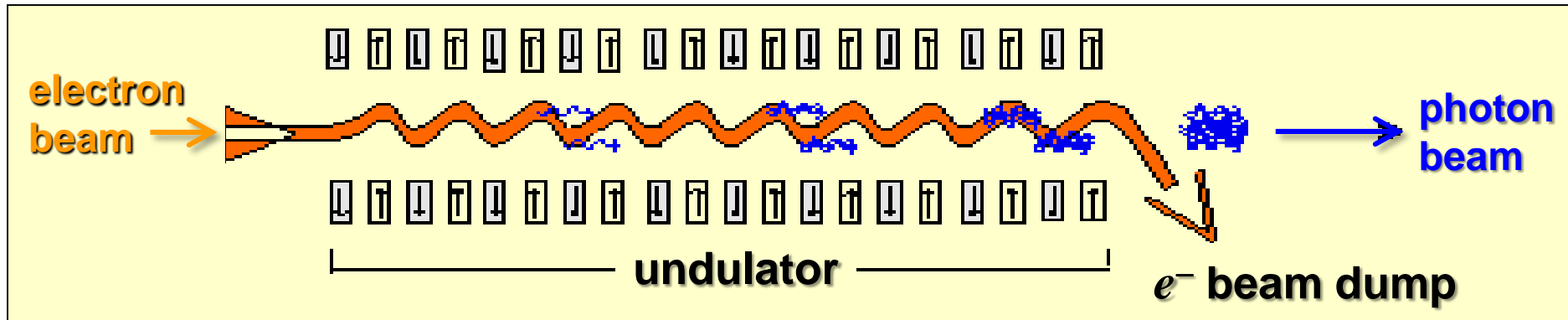
March 15, 2019

Coherence in particle and photon beams:
Past, Present, and Future Symposium

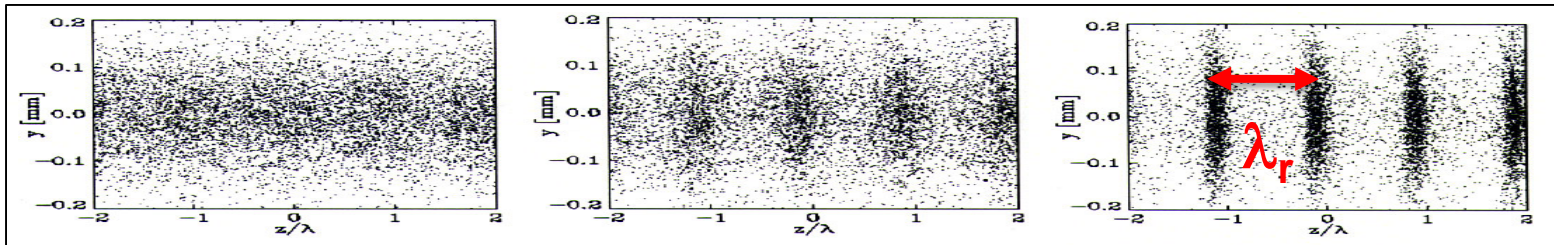


Free Electron Lasers

- Produced by **resonant interaction** of a relativistic electron beam with EM radiation in an undulator



$$\lambda_r = \frac{\lambda_u}{2\gamma^2} \left(1 + \frac{K^2}{2} \right)$$



- **Radiation intensity $\propto N^2$**
- **Tunable, Powerful, Coherent radiation sources**

Self-Amplified Spontaneous Emission (SASE)

- Initiated by electron shot noise (spontaneous emission) and amplified over a narrow frequency bandwidth $\sigma_\omega \sim \rho\omega_r$

K.-J. Kim, NIMA (1986), Wang & Yu, NIMA (1986)

$$\frac{dP}{d\omega} = g_A \left(\frac{dP_0}{d\omega} + g_S \frac{\rho\gamma_0 mc^2}{2\pi} \right) \exp \left(\frac{z}{L_G} - \frac{\Delta\omega^2}{2\sigma_\omega^2} \right)$$

↑ input power ↑

effective start-up noise power \approx undulator radiation over $2L_G$

- To determine the 3D effects including diffraction and finite beam size, one must solve the initial value problem in terms of a set of guided modes (first introduced by G. Moore)

Vlasov-Maxwell formalism

- The interaction between the electron beam and the FEL radiation can be described in the framework of the Vlasov-Maxwell equations.
- The e-beam is described in terms of a distribution function $F = F(\theta, \eta, \mathbf{x}, \mathbf{p}; z)$ in 6D-phase space. In view of the importance of stochastic effects such as shot noise, we use the [Klimontovich distribution](#):

$$F(\theta, \eta, \mathbf{x}, \mathbf{p}; z) = \frac{k_1}{n_e} \sum_{j=1}^{N_e} \delta[\theta - \theta_j(z)] \delta[\eta - \eta_j(z)] \\ \times \delta[\mathbf{x} - \mathbf{x}_j(z)] \delta[\mathbf{p} - \mathbf{p}_j(z)],$$

n_e : on-axis electron number density

- The distribution function is governed by the [Vlasov equation](#)

$$\frac{\partial F}{\partial z} + \frac{d\theta}{dz} \frac{\partial F}{\partial \theta} + \frac{d\eta}{dz} \frac{\partial F}{\partial \eta} + \frac{d\mathbf{x}}{dz} \cdot \frac{\partial F}{\partial \mathbf{x}} + \frac{d\mathbf{p}}{dz} \cdot \frac{\partial F}{\partial \mathbf{p}} = 0,$$

K.-J. Kim, PRL 57, 1871 (1986)

*K.-J. Kim, Z. Huang, R. Lindberg, Synchrotron Radiation and FELs (Cambridge Press, 2017)*⁴

Van Kampen's normal mode expansion

- After linearizing Vlasov Eq., we seek the **self-similar, guided eigenmodes** of the FEL. These are solutions of the form:

$$\Psi = \begin{bmatrix} a_\nu(\hat{\mathbf{x}}; \hat{z}) \\ f_\nu(\hat{\eta}, \hat{\mathbf{x}}, \hat{\mathbf{p}}, \hat{z}) \end{bmatrix} = e^{-i\mu_\ell \hat{z}} \begin{bmatrix} \mathcal{A}_\ell(\hat{\mathbf{x}}) \\ \mathcal{F}_\ell(\hat{\mathbf{x}}, \hat{\mathbf{p}}, \hat{\eta}) \end{bmatrix}$$

- They are characterized by a constant **growth rate** μ_l and a z-independent radiation/density mode profile A_l/F_l (**Optical guiding**)



- Substituting into the Vlasov-Maxwell (FEL) equations, we obtain two coupled relations for the growth rate and the mode amplitudes:

$$\begin{bmatrix} \mu_\ell \mathcal{A}_\ell + \left(-\frac{\Delta\nu}{2\rho} + \frac{1}{2} \hat{\nabla}_\perp^2 \right) \mathcal{A}_\ell + i \int d\hat{\mathbf{p}} d\hat{\eta} \mathcal{F}_\ell \\ \mu_\ell \mathcal{F}_\ell + i \mathcal{A}_\ell \frac{\partial \bar{f}_0}{\partial \hat{\eta}} + \left\{ -\nu \dot{\theta} + i \left(\hat{\mathbf{p}} \cdot \frac{\partial}{\partial \hat{\mathbf{x}}} - \hat{k}_\beta^2 \hat{\mathbf{x}} \cdot \frac{\partial}{\partial \hat{\mathbf{p}}} \right) \right\} \mathcal{F}_\ell \end{bmatrix} = 0.$$

3D solution

- Using Gaussian distributions, we obtain an explicit dispersion relation:

$$\left(\mu - \frac{\Delta\nu}{2\rho} + \frac{1}{2} \hat{\nabla}_{\perp}^2 \right) \mathcal{A}(\hat{x}) - \frac{1}{2\pi \hat{k}_{\beta}^2 \hat{\sigma}_x^2} \int_{-\infty}^0 d\tau \tau e^{-\hat{\sigma}_{\eta}^2 \tau^2 / 2 - i\mu\tau} \\ \times \int d\hat{p} \mathcal{A}[\hat{x}_+(\hat{x}, \hat{p}, \tau)] \exp \left[-\frac{1 + i\tau \hat{k}_{\beta}^2 \hat{\sigma}_x^2}{2\hat{k}_{\beta}^2 \hat{\sigma}_x^2} \left(\hat{p}^2 + \hat{k}_{\beta}^2 \hat{x}^2 \right) \right] = 0.$$

- There are four dimensionless parameters that affect the growth rate
(*L.-H. Yu, Krinsky, Gluckstern, Phys. Rev. Lett. 64, 1990*)
 - $\hat{\sigma}_x$ is a quantitative measure of the **diffraction effect**
 - $\hat{\sigma}_x \hat{k}_{\beta}$ is a measure of the **emittance effect**
 - $\hat{\sigma}_{\eta}$ represents the **energy spread effect**
 - $\Delta\nu/(2\rho)$ is scaled **frequency detuning**
- Ming Xie obtained a fitting formula that captures all these effects for FEL designs (1995)

Kwang-Je at APS since 1998

- Kwang-Je arrived at APS in 1998, I followed Kwang-Je to Chicago right after my Ph.D. from Stanford in May 1998.
- Our work was largely supported by an Argonne LDRD to do “**Comprehensive Analysis of SASE**”.
- In 1998-2002 we studied short-pulse effects, harmonic generation, 3D SASE start-up, FEL saturation, CSR microbunching instability...
- Together we published >20 journal publications and numerous conference papers, and went to some nice workshops too!

PHYSICAL REVIEW E VOLUME 62, NUMBER 5 NOVEMBER 2000

Three-dimensional analysis of harmonic generation in high-gain free-electron lasers

Zhirong Huang and Kwang-Je Kim
Advanced Photon Source, Argonne National Laboratory, Argonne, Illinois 60439

PHYSICAL REVIEW SPECIAL TOPICS - ACCELERATORS AND BEAMS, VOLUME 5, 074401 (2002)

Formulas for coherent synchrotron radiation microbunching in a bunch compressor chicane

Zhirong Huang* and Kwang-Je Kim
Argonne National Laboratory, Argonne, Illinois 60439

PHYSICAL REVIEW SPECIAL TOPICS - ACCELERATORS AND BEAMS 10, 034801 (2007)

Review of x-ray free-electron laser theory

Zhirong Huang
Stanford Linear Accelerator Center, Stanford, California 94309, USA

Kwang-Je Kim
Argonne National Laboratory, Argonne, Illinois 60439, USA

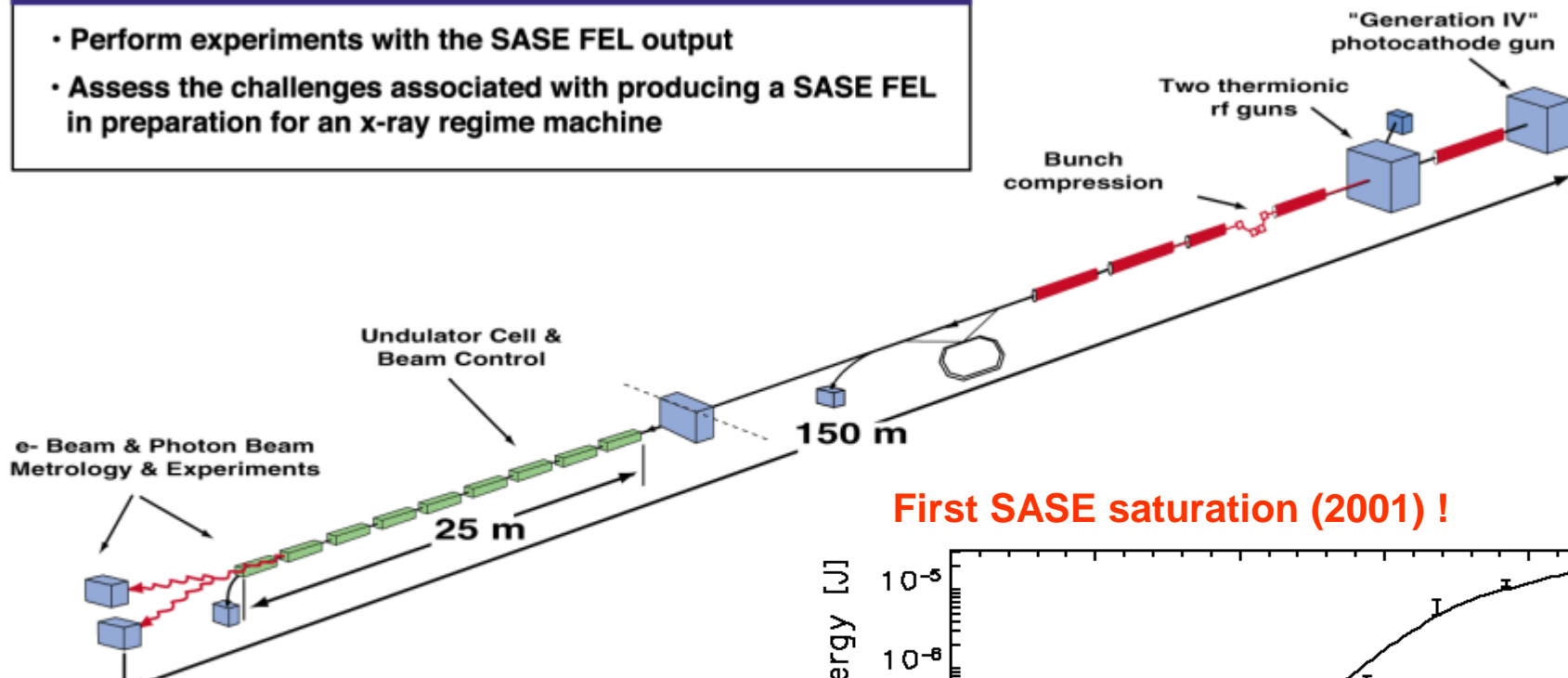


Sardinia beach (Italy 2002)

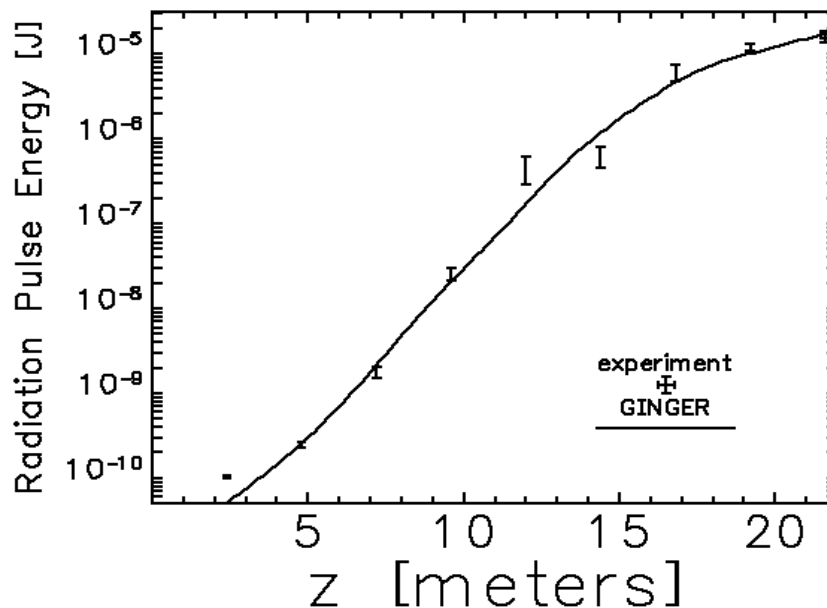
LOW-ENERGY UNDULATOR TEST LINE PARAMETERS

PROJECT GOALS

- Perform experiments with the SASE FEL output
- Assess the challenges associated with producing a SASE FEL in preparation for an x-ray regime machine



First SASE saturation (2001) !



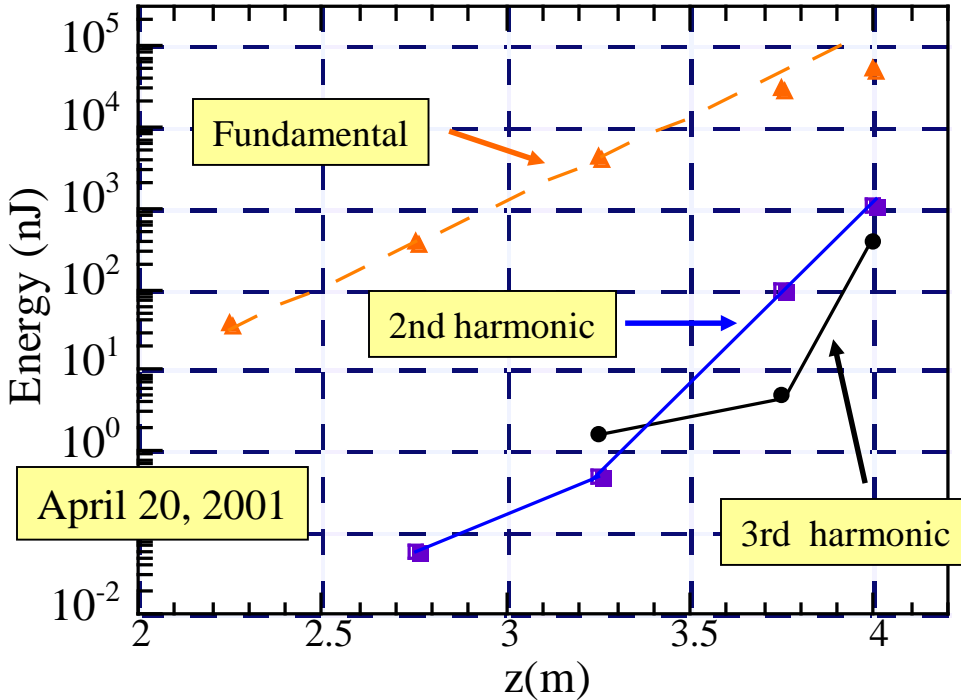
www.sciencemag.org SCIENCE VOL 292 15 JUNE 2001

Exponential Gain and Saturation of a Self-Amplified Spontaneous Emission Free-Electron Laser

S. V. Milton,^{1*} E. Gluskin,¹ N. D. Arnold,¹ C. Benson,¹ W. Berg,¹
 S. G. Biedron,^{1,2} M. Borland,¹ Y.-C. Chae,¹ R. J. Dejus,¹
 P. K. Den Hartog,¹ B. Deriy,¹ M. Erdmann,¹ Y. I. Eidelman,¹
 M. W. Hahne,¹ Z. Huang,¹ K.-J. Kim,¹ J. W. Lewellen,¹ Y. Li,¹
 A. H. Lumpkin,¹ O. Makarov,¹ E. R. Moog,¹ A. Nassiri,¹ V. Sajaev,¹
 R. Soliday,¹ B. J. Tieman,¹ E. M. Trakhtenberg,¹ G. Travish,¹
 I. B. Vasserman,¹ N. A. Vinokurov,³ X. J. Wang, G. Wiemerslage,¹
 B. X. Yang¹

Nonlinear Harmonic Radiation at VISA*

Nonlinear Harmonic Energy vs. Distance



Associated gain lengths

$$L_f = 19\text{cm}$$

$$L_2 = 9.8\text{cm} \quad \Rightarrow \quad L_n = L_g / n$$

$$L_3 = 6.0\text{cm}$$

Energy Comparison

Mode (n)	Wavelength (nm)	Energy (μJ)	% of E_1
1	845	52	
2	421	.93	1.8
3	280	.40	.77

Using the relation of 2nd and 3rd harmonic energies as given by Z. Huang and K.J.Kim

$$E_2 = \left(\frac{K}{\gamma k_u \sigma_x} \right)^2 \left(\frac{K_2}{K_3} \right)^2 \left(\frac{b_2}{b_3} \right)^2 E_3$$

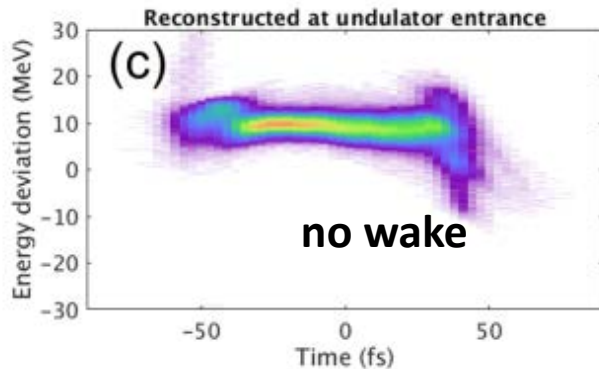
b -bunching parameters

K_n -Coupling coefficients

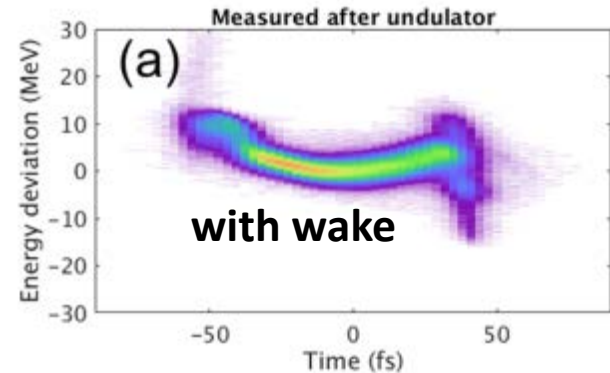
* A. Tremaine, XJ Wang et al., PRL (2002)

Onto LCLS

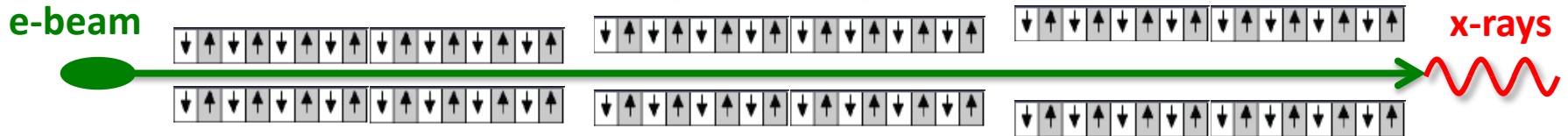
- I left Chicago for the Sunny California in late 2002.
- It was realized that undulator wakefield-induced energy loss is an important effect for the LCLS (5 mm gap for >100 m)



With XTCAV
(Y. Ding)



- Compensate the average energy loss by tapering undulator



- Tapered undulator keeps FEL resonance and increase power.
- But, undulator wakefield makes time-dependent energy loss and hence taper only works for the average loss.
- FEL resonance cannot be kept for every slice of the bunch.
- This led to FEL power degradation.

FEL with slowly varying beam and undulator parameters

➤ E-beam energy $\gamma_c(z)$, undulator parameter $K(z)$

➤ Initial resonant wavelength $\lambda_0 = \frac{2\pi}{k_0} = \frac{\lambda_u}{2\gamma_c(0)^2} \left[1 + \frac{K(0)^2}{2} \right]$

➤ Resonant energy $\gamma_r(z) = \sqrt{\frac{\lambda_u}{2\lambda_0} \left[1 + \frac{K(z)^2}{2} \right]}$

➤ Longitudinal motion is described by

$$\theta = (k_0 + k_u)z - ck_0 t^* \quad (\text{ponderomotive phase})$$

$$\eta = \frac{\gamma(z) - \gamma_c(z)}{\rho\gamma_c(0)} \quad (\text{normalized energy, change only due to FEL})$$

$$\frac{d\theta}{dz} = 2k_u \frac{\gamma(z) - \gamma_r(z)}{\gamma_c(0)} = 2k_u \rho \left[\eta + \frac{\gamma_c(z) - \gamma_r(z)}{\rho\gamma_c(0)} \right]$$

$$\frac{d\eta}{dz} \propto E \cos(\theta + \phi) \quad (E \text{ and } \phi \text{ are radiation field and phase})$$

WKB approximation

- Well-known technique in QM for slowly-varying potential
- FEL is characterized by ρ : the relative gain bandwidth is a few ρ , and radiation field gain length $\sim \lambda_u/(4\pi\rho)$

- Relative change in beam energy w.r.t resonant energy

$$\delta(z) = \frac{1}{\rho} \frac{\gamma_c(z) - \gamma_r(z)}{\gamma_c(0)} \quad \text{Normalized to } \rho$$

- Apply WKB technique if the relative energy change per field gain length is smaller than ρ , i.e.,

$$\left| \frac{d\delta}{d\tau} \right| < 1, \quad \tau = 2\rho k_u z = \frac{z}{\lambda_u/(4\pi\rho)}$$

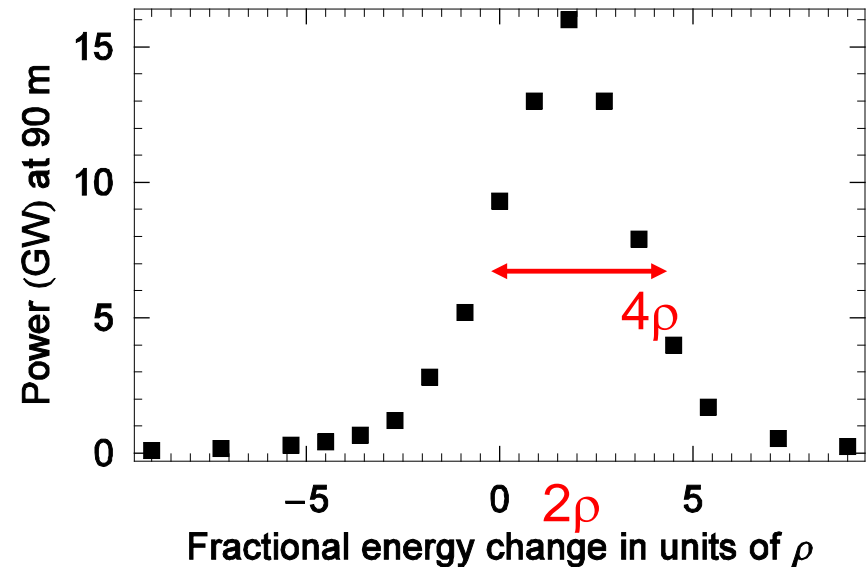
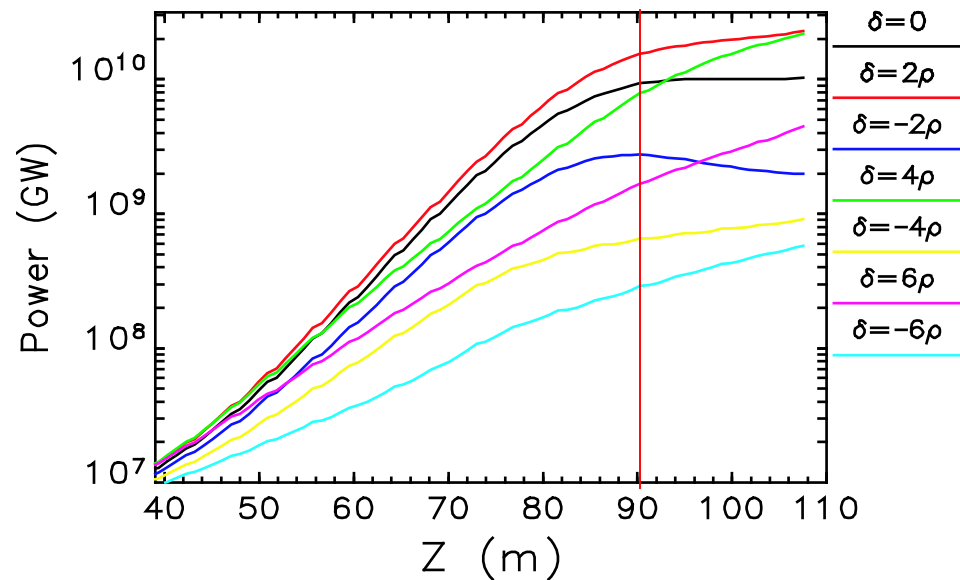
- We then extend the WKB analysis to 3D via Van Kampen's method of mode expansion.

Comparison w/ simulations

- Radiation power dependence on δ is a gaussian

$$P(\delta; z) = P_m(z) \exp \left[-\frac{1}{2} \left(\frac{\delta(z) - \delta_m(z)}{\sqrt{3}\sigma_\omega/\rho} \right)^2 \right]$$

- *GENESIS* simulation of LCLS power vs. δ ,
- ➔ Power enhancement ~ 2 when energy gain 2ρ at saturation

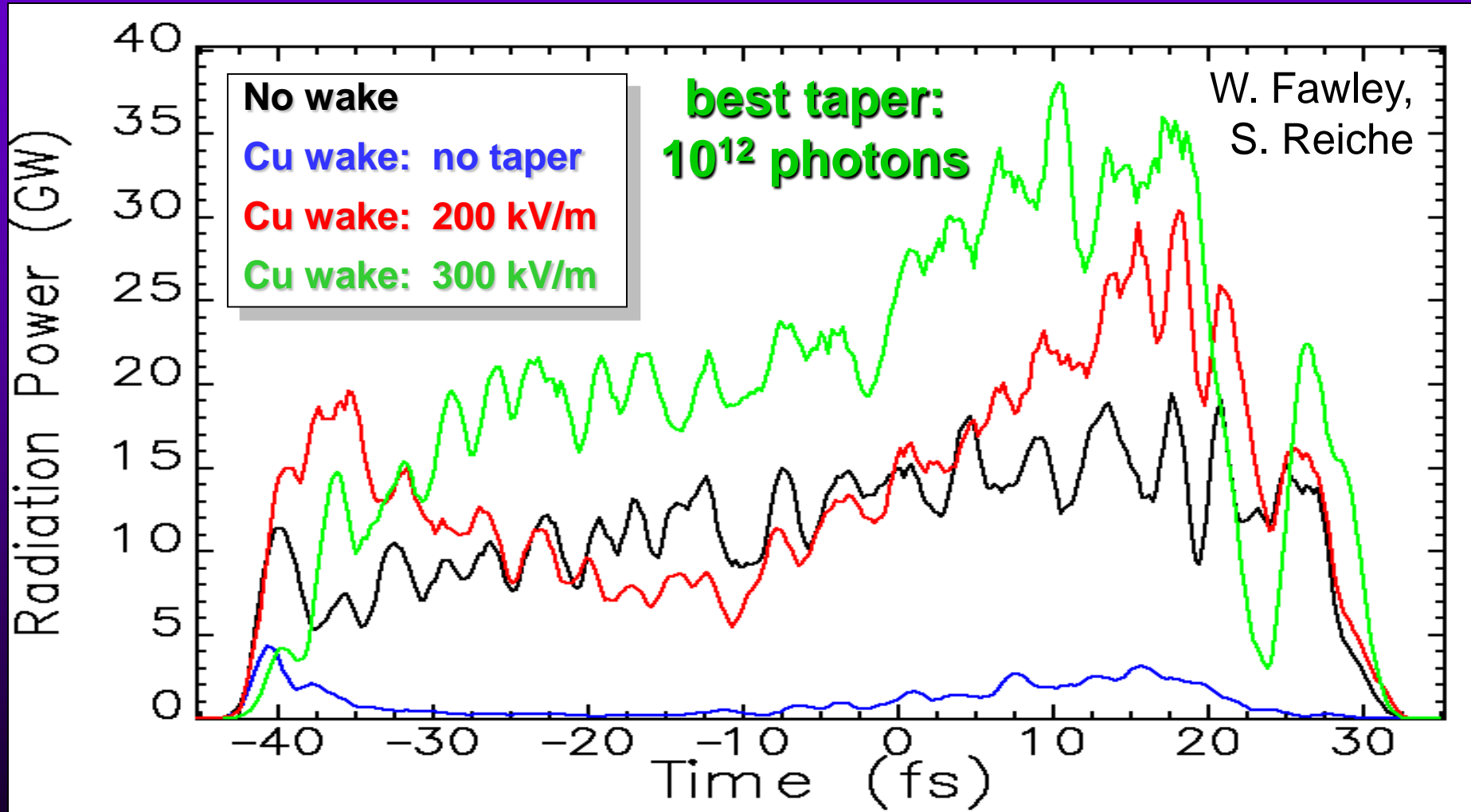


- Power vs. $\delta\rho$ has RMS $\sqrt{3}\sigma_\omega$

FWHM $4\sigma_\omega$ ($\sim 4\rho$ at saturation)

0.2-nC FEL Simulations with Taper

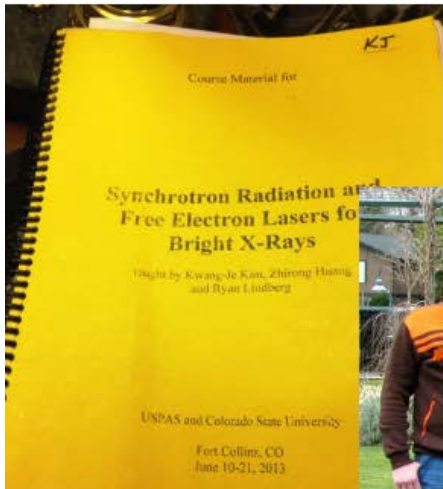
P. Emma's talk at PAC05



This study led to abandoning 1-nC LCLS

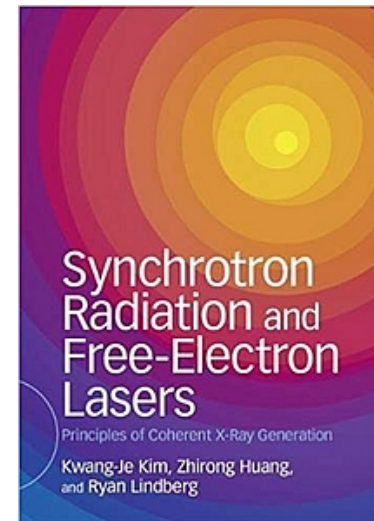
Teach FEL theory in USPAS

- KJK and I started the USPAS teaching in 2000, later joined force with Ryan. In total we have taught 8 USPAS sessions (+1 this coming summer).
- The lecture notes were steadily improved and became a textbook published by Cambridge Press in 2017.

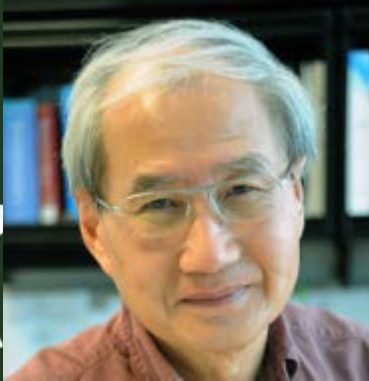


Lecture notes
2013

2008 USPAS at Santa Rosa, CA



- The book is translated into Chinese in 2018.
(Kwang-Je Kim -> 金光齐 -> Coherent radiation)



**MAY
THE
COHERENCE
BE WITH
YOU**

細推物理須行樂
何用浮名絆此身

杜甫 曲江二首

公元七五八年

*The law of Nature tells us to enjoy as we may.
Why spoil our joy by sheer vanity of life?*

**Poem by Fu Du (Tang dynasty, 758)
Calligraphy by T.-D. Lee (Nobel Laureate 1957)**